

It is evident that a quantitative theory for the diffusion of silver into glass must consider that the rate of diffusion depends on the glass, the melt and the atmosphere surrounding them. Such a theory is not yet available, and more experiments are probably needed in order to solve the problem. It is reasonable to believe, however, that the diffusion coefficients of other metals into Pyrex glass are of the same order of magnitude as that of silver¹⁵. Therefore the

exchange between a salt melt and a diffusion apparatus made of Pyrex glass should not be a serious problem. However, when the diffusion times are very long and when the glass surface exposed to the melt is very large (as in a fritted disc) special care should be taken.

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Self-Diffusion in Molten Silver Nitrate

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The self-diffusion coefficient of the Ag^+ -ion in molten AgNO_3 has been measured with the glass fibre paper strip technique in the temperature interval 220–336 °C. The result can be described by the equation

$$D = 3.2 \times 10^{-4} \exp\{-3400/(RT)\} \text{ cm}^2 \text{ s}^{-1}.$$

A comparison with the results obtained by other workers shows good agreement between the two sets of values. Possible sources of error are discussed.

Several different methods for self-diffusion measurements in molten salts have been used, e. g. different capillary methods¹, the shear cell method², and the fritted disc method^{2,3}. Another method, the paper strip technique, has been proposed by HONIG^{4,5}. Self-diffusion measurements in molten NaNO_3 , KNO_3 , and CsNO_3 showed good agreement between the results of this method and those obtained by previous workers⁶. In his original paper, HONIG discussed many aspects of using the paper strip method for diffusion measurements in molten AgNO_3 also, but he gave no experimental results. Since the relative merits of the different methods for diffusion measurements have been a matter of considerable discussion^{2–4,7,8}, we found that a closer investigation of the possibilities of the paper strip technique would be valuable. We decided to use it for in-

vestigation of the self-diffusion coefficient of the Ag^+ -ion in molten AgNO_3 . Such an investigation has been performed previously by DWORKIN, ESCUE and VAN ARTSDALEN⁹, using the open-ended capillary technique, and the exchange between an AgNO_3 melt and a borosilicate glass has been studied recently at this laboratory¹⁰. From a comparison between the results of these two investigations and the present results, conclusions could be drawn concerning the value of the paper strip technique.

Experimental

The experimental equipment was similar to that previously used by HONIG⁴. A strip, 15 cm long and 0.5 cm wide was cut from Whatman GF 81 (GF/A) borosilicate glass fibre paper. It was impregnated with molten AgNO_3 and placed on a Pyrex glass plate in a

¹ J. S. ANDERSON and K. SADDINGTON, J. Chem. Soc. London **1949**, Suppl. p. 381. — L. E. WALLIN, Z. Naturforsch. **17 a**, 191 [1962].

² S. DJORDJEVIC and G. J. HILLS, Trans. Faraday Soc. **56**, 269 [1960].

³ C.-A. SJÖBLOM and A. LUNDÉN, Z. Naturforsch. **18 a**, 942 [1963]. — C.-A. SJÖBLOM, Z. Naturforsch. **18 a**, 1247 [1963].

⁴ E. P. HONIG, Thesis, University of Amsterdam 1964.

⁵ J. A. A. KETELAAR and E. P. HONIG, J. Phys. Chem. **68**, 1596 [1964].

⁶ The comparison is made in ref. ⁴, page 60.

⁷ S. B. TRICKLEBANK, L. NANIS, and J. O'M. BOCKRIS, J. Phys. Chem. **68**, 58 [1964]. — J. O'M. BOCKRIS, S. YOSHIKAWA, and S. R. RICHARDS, J. Phys. Chem. **68**, 1838 [1964]. — L. NANIS, S. R. RICHARDS, and J. O'M. BOCKRIS, Rev. Sci. Instr. **36**, 673 [1965].

⁸ P. L. SPEDDING and R. MILLS, J. Electrochem. Soc. **112**, 594 [1965].

⁹ A. S. DWORKIN, R. B. ESCUE, and E. R. VAN ARTSDALEN, J. Phys. Chem. **64**, 872 [1960].

¹⁰ C.-A. SJÖBLOM and J. ANDERSSON, Z. Naturforsch. **21 a**, 274 [1966].



horizontal, wire-wound electrical furnace made from pyrophyllite¹¹. The temperature in the furnace was kept constant to $\pm 1^\circ\text{C}$ or better with a temperature controller. The temperature of the paper strip (and the melt) was measured with a thermocouple in contact with the glass plate. The thermocouple was connected to a Philips potentiometric recorder (measuring accuracy $\pm 0.5^\circ\text{C}$).

When temperature equilibrium had been established in the furnace, the radioactive tracer, $\text{Ag}^{110\text{m}}$ (obtained from Radiochemical Centre, Amersham, England) was brought on the paper. A little bit of radioactive salt was used in the original work by HONIG⁴, but such a technique does not lead to a true one-dimensional diffusion. In the present work, a very small paper strip, 5 mm \times 0.5 mm, impregnated with 50 μC $\text{Ag}^{110\text{m}}\text{NO}_3$, was placed across the centre of the big paper strip. A good contact between the small paper strip and the melt was obtained by pressing gently with a pair of tweezers. After a certain time (which was very short in comparison to the diffusion time) the small paper strip was removed and the diffusion proceeded. The duration of a diffusion anneal ranged from 6 to 18 hours. The paper strip stuck to the glass plate when the salt solidified after the run, thus no correction for the volume change of the salt was needed. The radioactivity distribution on the strip was measured with a scintillation counter, equipped with an 1 mm wide diaphragm made of lead. The self-diffusion coefficient can be calculated from this distribution and the diffusion time in a way described elsewhere⁴.

In order to see if any exchange between the salt melt and the glass took place during the diffusion anneal, the radioactivity measurement was repeated when the AgNO_3 had been washed away from the strip. We found a certain amount of $\text{Ag}^{110\text{m}}$ activity in the glass, the effect of which will be discussed below.

Results and Discussion

The experimental results (corrected for the influence of the glass radioactivity) are given in the table, and they can be described by the following ARRHENIUS' equation, obtained by "least squares" fitting¹²:

$$D = 3.2 \times 10^{-4} \exp\left(-\frac{3400 \pm 320}{RT}\right) (\text{cm}^2 \text{ s}^{-1})$$

where R is expressed in cal mole⁻¹ degree⁻¹ and T in degrees Kelvin.

The present value of the activation energy, $Q = 3.4$ kcal/mole, is within experimental error equal to the value obtained by DWORKIN, ESCUE, and VAN ARTSDALEN⁹, using the conventional, open-ended

T $^\circ\text{C}$	$D_{\text{Ag}^+} \times 10^5$ $\text{cm}^2 \text{ s}^{-1}$
220	0.95
236	1.04
249	1.07
268	1.21
274	1.23
278	1.24
288	1.45
306	1.43
320	1.59
336	1.96

Table 1. Self-diffusion coefficient of the Ag^+ -ion in molten silver nitrate, measured with the glass fibre paper strip technique.

capillary technique¹. The present D -values and the previous ones are in excellent agreement with each other at 220°C , and they differ slightly at 275°C . Thus the results of these two independent methods are seen to confirm each other. The reproducibility of the paper strip technique is not quite as good as that of the capillary method, if $\ln(D \times 10^6)$ is plotted against $10^3/T$, the standard deviation of the experimental points is 2.5%.

The cause of the somewhat poorer reproducibility of the paper-strip method is twofold: It is very difficult to impregnate the paper strip uniformly enough. A difference in "weight per cm" of 3% between different parts of the strip, which is considered by HONIG to be acceptable, might give rise to a difference in D of about 8%. Moreover, the exchange of silver between the melt and the glass affects the value of the experimental self-diffusion coefficient. After the diffusion run, about 3% of the total $\text{Ag}^{110\text{m}}$ radioactivity is found to remain in the centre of the glass fibre paper strip, cf. ref.⁴, page 101. The Ag^+ -ions react further with the glass when they have penetrated into it, which means that the remaining radioactivity inside the glass corresponds to the maximum radioactivity of the outside melt during the experiment. Thus it is found, when the final radioactivity distribution on the strip is measured, that the nearer to the maximum of the distribution curve, the larger a fraction of the measured activity is contributed by the radioactivity trapped in the glass. This effect makes the experimental self-diffusion coefficients too low, and a mean correction of about 8% had to be applied in the present work. This was determined in the separate measurement of the radioactivity distribution in the glass. It is very difficult to obtain good counting statistics in this

¹¹ Obtained under the name of Wonderstone from Carters (Merchants) Ltd., London.

¹² A. HALD, Statistical Theory with Engineering Applications, John Wiley & Sons, New York 1952, p. 522. All stated errors in this paper are standard deviations.

measurement, which makes the correction factor rather uncertain.

The influence of end-effects (Δl -effects) on the results of the capillary method has recently been discussed by several authors. BOCKRIS et al.⁷ claim that it is impossible to avoid Δl -errors even if the flow rate past the capillary mouth is carefully chosen. According to SPEDDING and MILLS⁸, on the other hand, the Δl error will be negligible for certain flow rates. Since the investigation by DWORKIN et al. was performed before the importance of these end effects was fully appreciated, it is reasonable to suspect that end errors in their capillary method might have affected their results. However, since our values and their values of D_{Ag^+} in molten $AgNO_3$ are in good agreement, we find that DWORKIN et al. have succeeded in avoiding Δl errors (at least in their measurements of D_{Ag^+}).

The above results show, that the paper strip technique can be used for self-diffusion measurements

in molten salts if a correction for the exchange of metal ions between the glass and the melt is applied when necessary. The glass-melt interface is very large, and this fact, in conjunction with the comparatively long diffusion times that are needed, enhances the exchange. (A somewhat similar situation is encountered in the porous frit technique^{2,3}, but the corresponding problem is by no means so serious, since the diffusion times are of the order of 30 minutes only.) Due to the difficulty to determine the correction factor accurately, the paper strip method is not capable of the highest precision. It is, on the other hand, very simple to use, and it is able to give valuable independent checks of the results of other methods.

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Thermoelectric Power of Molten and Solid Silver Sulphate

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The thermocell $Ag(T_1)/Ag_2SO_4/Ag(T_2)$ has been studied over the temperature range 420 to 750 °C. The results can be described by the following linear relations¹

$$-\varepsilon_1 = (0.31 \pm 0.01) \text{ mV/}^\circ\text{C} \quad (657-750^\circ\text{C})$$

$$-\varepsilon_s = [(0.687 \pm 0.011) - (0.472 \pm 0.019) \cdot 10^{-3} \cdot t_m] \text{ mV/}^\circ\text{C} \quad (420-657^\circ\text{C})$$

ε_1 and ε_s are the SEEBECK-coefficients of molten silver sulphate and the α -modification, respectively, and t_m is the mean temperature of the salt column in °C.

The standard deviation of ε_s was $4.3 \cdot 10^{-3} \text{ mV/}^\circ\text{C}$.

The transported entropy of the silver ion and of the sulphate ion have been calculated. A discontinuous change of the emf and the transported entropies was found at the melting point.

A number of measurements of the thermoelectric power of solid and molten ionic salts has been reported^{2,3}.

Between 412 and 657 °C silver sulphate forms a high temperature modification with a hexagonal structure⁴. This modification is characterised by a high cation mobility. The electrical conductivity increases strongly at the β — α transition and the conductivity of the α -modification is close to that of the melt. Similar high temperature modifications have been found in Li_2SO_4 and K_2SO_4 .

We report here on the determination of the thermoelectric power of the system $Ag(T_1)/Ag_2SO_4/Ag(T_2)$ in the temperature range 420 to 750 °C.

Experimental

Reagent grade silver sulphate (Hopkin & Williams or May & Baker) has been used in all experiments without further purification. No detectable difference between the results from two lots could be found, and both lots were therefore considered to be of satisfactory purity.

¹ All stated errors in this paper are standard deviations.

² A. KVIST and A. LUNDÉN, Z. Naturforschg. **19 a**, 1058 [1964] and references there.

³ A. KVIST, Z. Naturforschg. **19 a**, 1159 [1964].

⁴ H. A. ØYE, Acta Chem. Scand. **18**, 361 [1964]; Thesis Trondheim 1963.